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## STRESSES IN DEEP BEAMS

By Li Chow, Harry D. Conway, and  
George Winter, M. ASCE

STRUCTURAL DIVISION

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<i>Technical Division</i>	<i>Proceedings-Separate Number</i>
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City Planning.....	58, 60, 62, 64, 93, 94, 99, 101, 104, 105, 115, 131 (Discussion: D-16, D-23, D-43, D-60, D-62, D-65)
Construction.....	43, 50, 55, 71, 92, 94, 103, 108, 109, 113, 117, 121, 126, 130, 132 (Discussion: D-3, D-8, D-17, D-23, D-36, D-40, D-71)
Engineering Mechanics.....	41, 49, 51, 54, 56, 59, 61, 88, 89, 96, 116, 122, 124, 125, 126, 127, 128, 129 (Discussion: D-5, D-XXIII, D-XXV, D-18, D-24, D-33, D-34, D-49, D-54, D-61)
Highway.....	43, 44, 48, 58, 70, 100, 105, 108, 113, 120, 121, 130 (Discussion: D-XXVIII, D-7, D-13, D-16, D-23, D-60)
Hydraulics.....	50, 55, 56, 57, 70, 71, 78, 79, 80, 83, 86, 92, 96, 106, 107, 110, 111, 112, 113, 116, 120, 123, 130 (Discussion: D-XXVII, D-9, D-11, D-19, D-28, D-29, D-56, D-70, D-71, D-79)
Irrigation and Drainage.....	46, 47, 48, 55, 56, 57, 67, 70, 71, 87, 88, 90, 91, 96, 97, 98, 99, 102, 106, 109, 110, 111, 112, 114, 117, 118, 120, 129, 130 (Discussion: D-XXIII, D-3, D-7, D-11, D-17, D-19, D-25-K, D-29, D-30, D-38, D-40, D-44, D-47, D-57, D-70, D-71, D-80)
Power.....	48, 55, 56, 69, 71, 88, 96, 103, 106, 109, 110, 117, 118, 120, 129, 130 (Discussion: D-XXIII, D-2, D-3, D-7, D-11, D-17, D-19, D-25-K, D-30, D-38, D-40, D-44, D-70, D-71, D-79)
Sanitary Engineering.....	55, 56, 87, 91, 96, 106, 111, 118, 130 (Discussion: D-10, D-29, D-37, D-56, D-60, D-70, D-79, D-80)
Soil Mechanics and Foundations.....	43, 44, 48, 94, 102, 103, 106, 108, 109, 115, 130 (Discussion: D-4, D-XXVIII, D-7, D-43, D-44, D-56)
Structural.....	42, 49, 51, 53, 54, 59, 61, 66, 89, 100, 103, 109, 113, 116, 117, 119, 121, 122, 123, 124, 125, 126, 127, 128, 129, 132 (Discussion: D-32, D-33, D-34, D-37, D-39, D-42, D-49, D-51, D-53, D-54, D-59, D-61, D-66, D-72)
Surveying and Mapping.....	50, 52, 55, 60, 63, 65, 68, 121 (Discussion: D-60, D-65)
Waterways.....	41, 44, 45, 50, 56, 57, 70, 71, 96, 107, 112, 113, 115, 120, 123, 130 (Discussion: D-8, D-9, D-19, D-27, D-28, D-56, D-70, D-71, D-79, D-80)

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PAPERS

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STRESSES IN DEEP BEAMS

BY LI CHOW,<sup>1</sup> HARRY D. CONWAY,<sup>2</sup> AND GEORGE  
WINTER,<sup>3</sup> M. ASCE

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SYNOPSIS

Beams whose depths are comparable to their spans are used in a variety of structures. The distribution of bending and shear stresses in such deep beams departs radically from that given by the ordinary, simple formulas for shallow members. Information on the stresses in continuous, deep beams is available elsewhere, and corresponding information for single-span beams is presented in this paper. Five cases of loading are studied, and, for four of these cases, three different span-to-depth ratios are examined. Distributions and magnitudes of bending and shear stresses are given in graphical and tabular form suitable for direct use in design. Although this information is directly applicable to structures made of homogeneous material, such as steel, their use in connection with reinforced concrete requires some special considerations that are briefly outlined.

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INTRODUCTION

A deep beam may be defined as one whose depth is comparable to its span. Beams of this type, both in steel and in reinforced concrete, often arise in the construction of bins, hoppers, or similar structures, as well as in more ordinary construction in foundation walls or in cases in which walls are supported on individual columns or footings. The horizontal or vertical diaphragms used to transmit wind forces in buildings (floors or walls) are frequently of such dimensions as to represent deep beams. In reinforced concrete hipped-plate construction<sup>4</sup> the plates of the structure proper or the supporting diaphragms often fall into this category.

In all these cases, design based on the ordinary, straight-line distribution of bending stresses in shallow beams may be seriously in error, since the simple

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NOTE.—Written comments are invited for publication; the last discussion should be submitted by November, 1952.

<sup>1</sup> Graduate student at Cornell Univ., Ithaca, N. Y.

<sup>2</sup> Prof. of Applied Mech., Cornell Univ., Ithaca, N. Y.

<sup>3</sup> Prof. and Head, Dept. of Structural Eng., Cornell Univ., Ithaca, N. Y.

<sup>4</sup> "Hipped Plate Construction," by G. Winter and M. Pei, *Journal, Am. Concrete Inst.*, Vol. 18, No. 5, 1947, pp. 505-531.

theory of flexure (according to Navier's hypothesis) takes no account of the effect of the normal pressures on the top and bottom edges of the beam caused by the loads and reactions. The effect of these normal pressures on the stress distribution in deep beams is such that the distribution of bending stresses on vertical sections is not linear and the distribution of shear stresses is not parabolic. Consequently, a transverse section which is plane before bending does not remain approximately plane after bending and the neutral axis does not usually lie at the mid-depth, its position being variable in a span-wise direction.

Since information on the stresses in continuous deep beams have already been made available<sup>5</sup> as a result of the work of F. Dischinger,<sup>6</sup> the purpose of this paper is to present the stresses in single-span beams under various types of loading and to indicate how they differ in distribution and magnitude from those predicted by the ordinary beam formulas. It is assumed that the beam material is homogeneous and isotropic. The beams are analyzed as problems in plane stress, using the finite difference method to solve the differential equation for the stress function. The complete details of these solutions have been given by Mr. Chow,<sup>7</sup> while the finite difference method itself has been fully described and discussed in a previous paper by two of the present authors.<sup>8</sup> For this reason, the method proper will not be described in any detail although the results obtained from it are given in full. Although most of these results are new, some obtained previously by H. Bay<sup>9</sup> have been included for completeness and are so indicated on the corresponding distribution charts.

#### FORMULAS OF THE FINITE DIFFERENCE METHOD

The area enclosed by the four edges of a deep beam is designated as the  $x$ - $y$  plane, and this plane is divided into a network of equal divisions by lines parallel to the edges;  $h$  and  $k$  are the lengths of divisions in the  $x$  and  $y$  directions, respectively. Furthermore,  $Z_{x,y}$  denotes the ordinate at each net-point  $(x,y)$  to a curved surface representing G. B. Airy's stress function  $z = f(x,y)$ . Then, the biharmonic differential equation that Airy's stress function must satisfy is equivalent to the following linear equation in terms of  $Z$ :

$$\begin{aligned} Z_{x,y} \left\{ 6 \left( \lambda + \frac{1}{\lambda} \right) + 8 \right\} - 4 \left\{ (1 + \lambda) (Z_{x-h,y} + Z_{x+h,y}) \right. \\ \left. + \left( 1 + \frac{1}{\lambda} \right) (Z_{x,y-k} + Z_{x,y+k}) \right\} + 2(Z_{x-h,y-k} + Z_{x-h,y+k} + Z_{x+h,y-k} \\ + Z_{x+h,y+k}) + \lambda(Z_{x-2h,y} + Z_{x+2h,y}) + \frac{1}{\lambda}(Z_{x,y-2k} + Z_{x,y+2k}) = 0 \dots (1) \end{aligned}$$

in which  $\lambda = (k/h)^2$ .

<sup>5</sup> "Design of Deep Girders," Pamphlet No. ST 66, Concrete Information, Structural Bureau, Portland Cement Assn., Chicago, Ill.

<sup>6</sup> "Beitrag zur Theorie der Halbscheibe und des wandartigen Trägers," by F. Dischinger, *Publications, International Assn. for Bridge and Structural Eng.*, Zurich, Switzerland, Vol. 1, 1932, pp. 69-93.

<sup>7</sup> "Stresses in Deep Beams," by Li Chow, thesis presented to Cornell University, at Ithaca, N. Y., in 1951, in partial fulfillment of the requirements for the degree of Doctor of Philosophy.

<sup>8</sup> "Analysis of Deep Beams," by H. D. Conway, L. Chow, and G. W. Morgan, *Journal of Applied Mechanics*, Am. Soc. of Mech. Engrs., Vol. 18, No. 2, June, 1951, pp. 163-172.

<sup>9</sup> "Über den Spannungszustand in hohen Trägern und die Bewehrung von Eisenbetontragwänden," by H. Bay, Stuttgart, Germany, 1931, p. 64.

The unknown  $Z$ -values are determined by solving the set of simultaneous linear equations obtained from the application of Eq. 1 to each net-point. The normal stress  $\sigma$  at any point  $(x, y)$  can then be calculated by means of the following formula,

$$\sigma_x = \frac{\partial^2 Z}{\partial y^2} = \frac{Z_{x, y-k} - 2Z_{x, y} + Z_{x, y+k}}{k^2} \dots \dots \dots (2a)$$

$$\sigma_y = \frac{\partial^2 Z}{\partial x^2} = \frac{Z_{x-h, y} - 2Z_{x, y} + Z_{x+h, y}}{h^2} \dots \dots \dots (2b)$$

The shear stress  $\tau$  at the point is

$$\tau_{xy} = -\frac{\partial^2 Z}{\partial x \partial y} = \frac{(Z_{x-h/2, y+k/2} + Z_{x+h/2, y-k/2}) - (Z_{x-h/2, y-k/2} + Z_{x+h/2, y+k/2})}{hk} \dots (3)$$

It should be noticed that Eqs. 2(a) and 2(b) give the normal stresses on the sections along the lines of division of the network, whereas Eq. 3 gives the shear stresses on sections midway between the lines of division, due to the presence of  $h/2$  and  $k/2$  in the subscripts of  $Z$ .

#### TYPES OF LOADING

Five types of loading have been analyzed by the finite difference method and are indicated in Fig. 1. These loadings were chosen to represent, in a general manner, the most common types occurring in practice. Each of the first four types (Fig. 1(a) to Fig. 1(d)) were investigated for three values of the height-to-span ratio, namely,  $H/L = \frac{1}{2}$ , 1, and 2. The results presented herein utilizing the loading of Fig. 1(a) (with  $H/L = \frac{1}{2}$  and 1) and Fig. 1(c) (with  $H/L = 1$  and 2) are taken from the analyses made by Mr. Bay, who also used the finite difference method.

For the loading of Fig 1(e), the central portion of the beam is subject to pure bending, and the simple flexure theory predicts zero shear stresses throughout this portion. However, this is not true for deep beams, except for the midspan section in which the shear stresses are zero by symmetry. The case of  $H/L = 1$  was, therefore, analyzed to show this deviation from the usual assumption.

For all the cases presented in this paper,  $h = L/6$  and  $k = H/6$  were used in constructing the network of divisions. However, Mr. Bay used  $h = L/4$  and  $k = H/4$  for the case of  $H/L = \frac{1}{2}$  under the loading of Fig. 1(a) and  $h = L/6$  and  $k = H/5$  for the case of  $H/L = 2$  under the loading of Fig. 1(c). A somewhat better accuracy was, thereby, achieved in the writers' solutions than in those of Mr. Bay.

#### RESULTS AND DISCUSSION

Figs. 2(a), 3(a), 4(a), and 5(a) show the bending stress distributions at the midspan section of a beam under the loadings of Fig. 1(a) to Fig. 1(d). The stress shown is unit stress, being in terms of beam width,  $b$ . The midspan section was chosen for consideration since it is the section of maximum bending

moment. It is seen that, for all types of loading, when  $H/L = \frac{1}{2}$  the stress curve agrees reasonably well with the linear distribution of the simple flexure theory. The deviation of the curves from the linear distribution becomes increasingly pronounced as the height-to-span ratio increases. The maximum

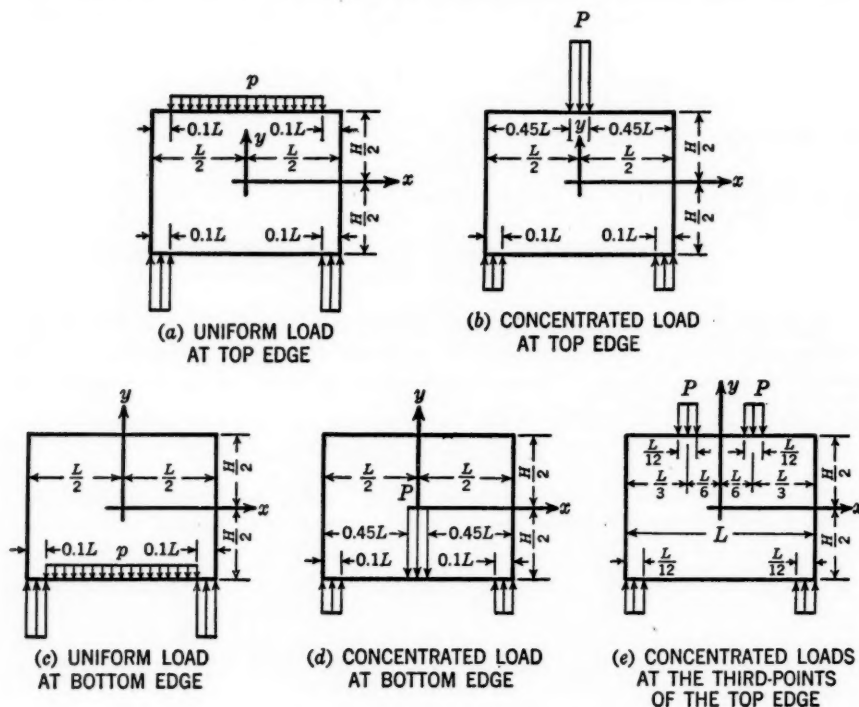


FIG. 1.—VARIOUS TYPES OF LOADING

bending stress for values of  $H/L \geq 1$  is considerably greater than that predicted by the linear theory, but the stress decreases very rapidly with increasing distance from the edge of maximum stress. There are three points of zero bending stress when  $H/L = 2$  with the load acting along the top edge, although the upper half is almost free from bending stress when the load acts along the bottom edge.

Figs. 2(b), 3(b), 4(b), and 5(b) give the shear stress distributions at the section  $x = L/4$  under the loading shown in Fig. 1(a) to Fig. 1(d). The total shear force at this transverse section is a maximum for the cases shown in Fig. 1(b) and Fig. 1(d), but it is less than the maximum value that occurs at the section  $x = 0.4L$ , for the cases shown in Fig. 1(a) and Fig. 1(c). Deviations from the simple theory are always present for sections near the loads and reactions, no matter what value the height-to-span ratio may be. These variations are the result of the local influence of the loading and occur also in shallow beams. For this reason, it is preferable to consider a section that is at some distance from the supports for investigating the validity of the simple theory. Moreover, the



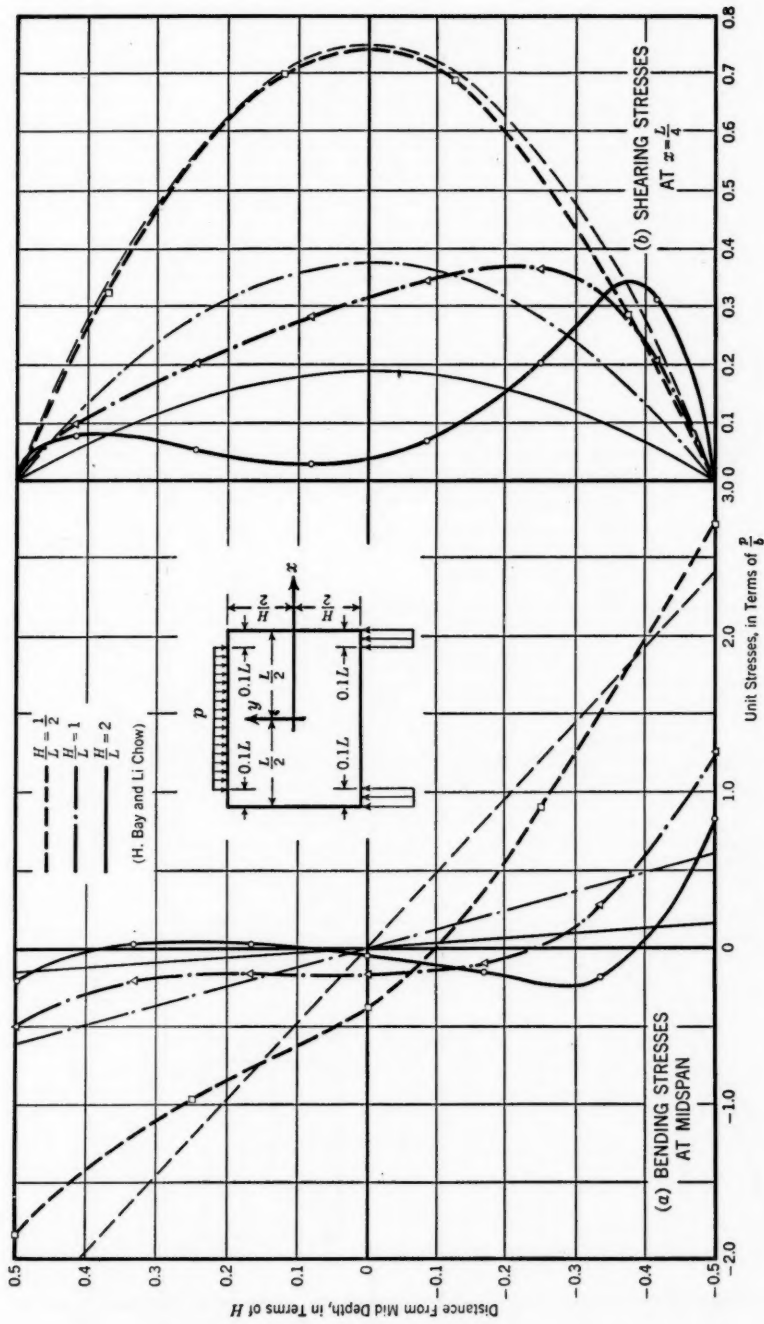


FIG. 2.—STRESSES IN BEAMS HAVING UNIFORM LOADING AT THE TOP EDGE

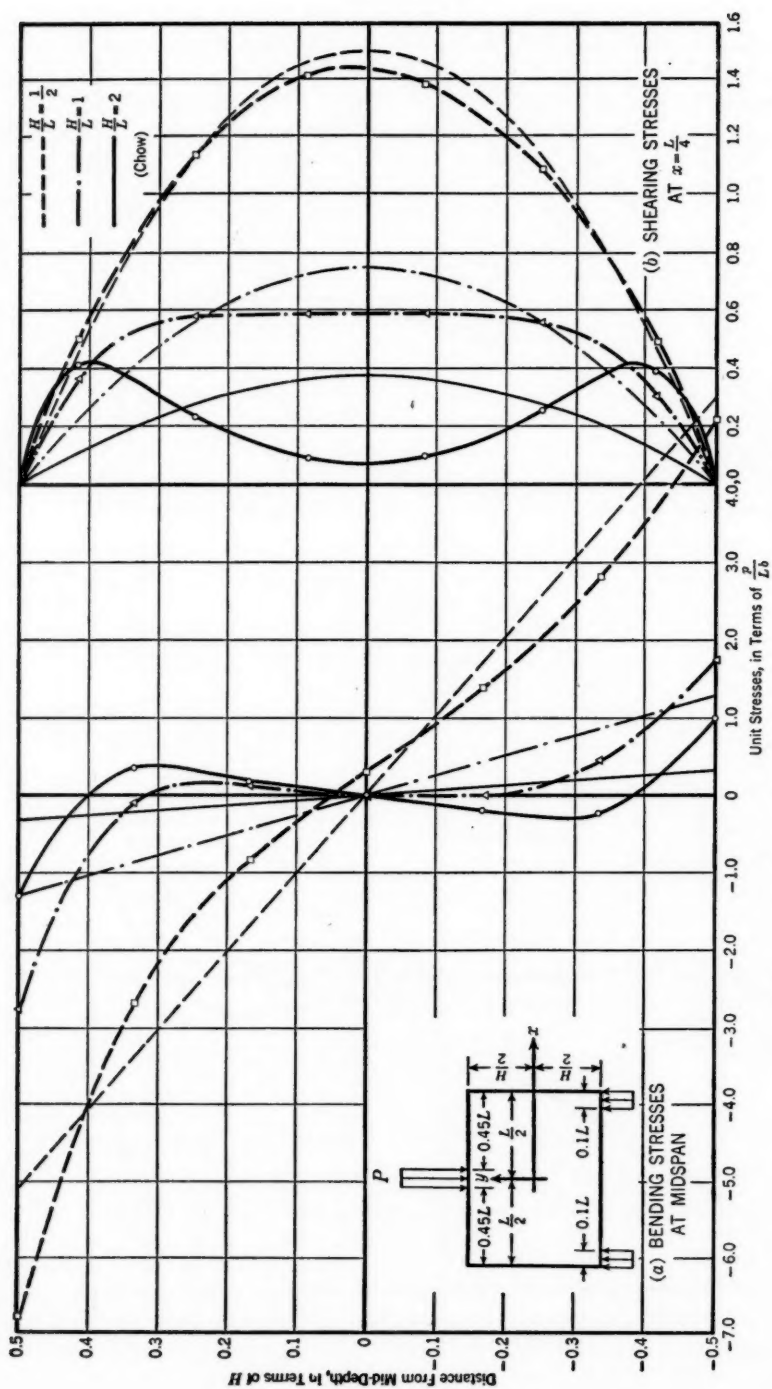


FIG. 3.—STRESSES IN BEAMS HAVING CONCENTRATED LOADING AT THE TOP EDGE



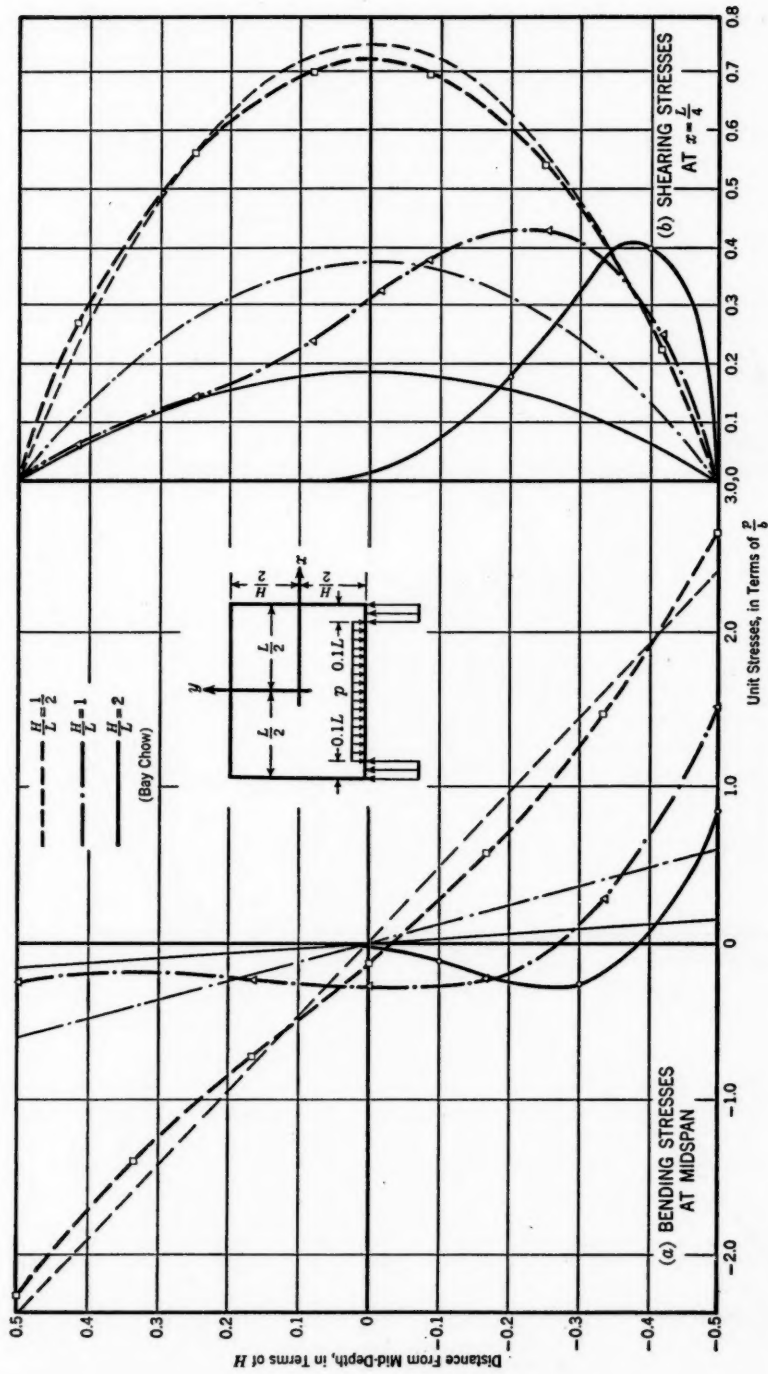


FIG. 4.—STRESSES IN BEAMS HAVING UNIFORM LOADING AT THE BOTTOM EDGE

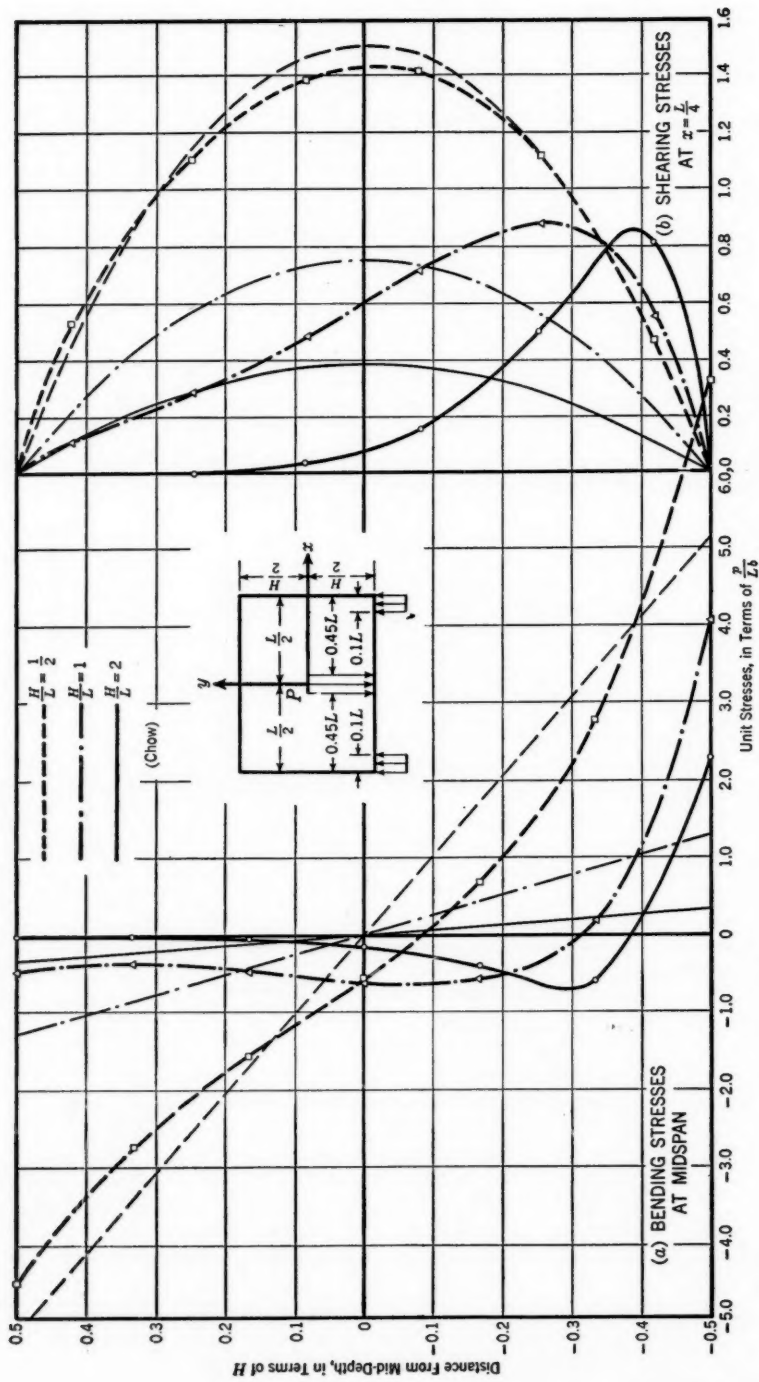


FIG. 5.—STRESSES IN BEAMS HAVING CONCENTRATED LOADING AT THE BOTTOM EDGE

stress values obtained from the finite difference method are usually inaccurate for regions in which abrupt changes of stress and boundary conditions take place. At sections in which the total shear force varies suddenly, the shear stresses computed by this method are of little value.

It is seen from Figs. 2(b), 3(b), 4(b), and 5(b) that the shear stress curves differ radically from the conventional parabolic form when the height-to-span ratio is great. Moreover, the shear stress vanishes in the upper half of the depth, for the cases of  $H/L = 2$  with loads acting at the bottom edge. It may also be observed that the areas under these shear stress curves correspond to forces that are somewhat smaller than the actual shear force at the section. This is caused by the inaccuracy of the finite difference method. For this reason, Figs. 2(b), 3(b), 4(b), and 5(b) may be taken as indicating only the probable variations in shear stress and, in so far as actual magnitudes are concerned, these figures must be used with some discretion. However, the distributions of bending stresses given by the finite difference method are known to be quite accurate, and these stresses are probably the most important from a practical point of view.

Fig. 6(a) shows the bending stress distribution at sections  $x = 0$  and  $x = L/6$  under the loading of Fig. 1(e), with  $H/L = 1$ . These sections have nearly equal bending moments ( $0.292 PL$  for  $x = 0$  and  $0.281 PL$  for  $x = L/6$ ), but their stress curves are quite different. The maximum compressive bending stress at the upper edge for the section ( $x = L/6$ ) is about twice that for the section ( $x = 0$ ), although the latter section has a bending moment that is slightly greater. This indicates the presence of large localized stresses in the vicinity of concentrated loads. It may also be pointed out from an examination of Fig. 6(a) that the shapes of the bending stress curves for two adjacent transverse sections are not similar to each other and their neutral axes do not necessarily coincide.

Fig. 6(b) shows the shear stress distributions at the sections  $x = L/12$ ,  $L/4$ , and  $5L/12$  for the loading of Fig. 1(e), with  $H/L = 1$ . The peaks of the stress curves are near the top edge for sections near the central portion of the span, but these peaks are near the bottom edge for sections near the supports. Although the total shear force is zero at the section ( $x = L/12$ ), the shear stresses do not vanish in this section. It has been shown<sup>7</sup> that in the neighborhood of a concentrated load, localized shear stresses are induced near the discontinuity of the large normal pressure. Thus, localized shear stresses are produced at the section  $x = L/12$ , in the vicinity of the concentrated load, and shear stresses of opposite sign must also be induced in other parts of the same section, since the total shear force is zero.

The inaccuracy of the shear stress values in Fig. 6(b) is more pronounced than that in Figs. 2(b), 3(b), 4(b), and 5(b), since each of the three sections considered is located close to a point of sudden change of shear force. As may be observed from Fig. 6(b), the algebraic sum of the areas under the curve for the section  $x = L/12$  is not zero and the areas under the curves for the other two sections correspond to forces that are appreciably less than the total shearing force  $P$  indicated by the area under the parabola.

All the stress values given in Figs. 2 to 6 are either in terms of  $p/b$ , in which  $p$  is the uniform load per unit length of the beam and  $b$  is the width of the beam, or in terms of  $P/(Lb)$ , in which  $P$  is the concentrated load and  $L$  is the length of the beam, according to the type of loading indicated in each of the figures.

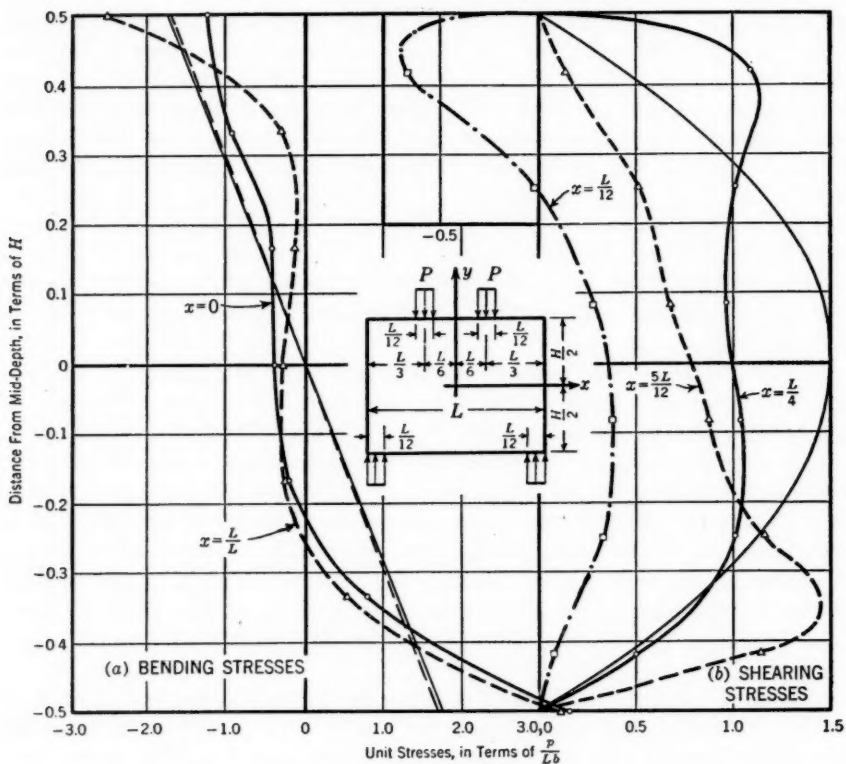


FIG. 6.—STRESSES FOR BEAMS WITH  $\frac{H}{L} = 1$  AND HAVING CONCENTRATED LOADS AT THE THIRD POINTS OF THE TOP EDGE

By means of the principle of superposition, it is possible to compute the stresses for any combination of the given loadings, such as for uniform loading with a concentrated load at the center of the span. For example, let it be required to find the extreme fiber stresses at the midspan section of a single-span deep beam for which the following data are given: Length of support for reaction or concentrated load = 3 ft; span ( $L$ ) = 30 ft; height of beam ( $H$ ) = 30 ft; width of beam ( $b$ ) = 1.25 ft; uniform load along top edge = 30,000 lb per linear ft = 2,500 lb per in.; and concentrated load at center of top edge = 100,000 lb. Since  $H/L = 1$ , the stress coefficients selected from Figs. 2(a) and 3(a) are  $-0.50$  and  $+1.24$  for the uniform loading and  $-2.74$  and  $+1.74$  for the concentrated load. The plus sign indicates tension and the minus sign, compression. The maximum and minimum bending stresses in pounds per square

inch for the combined loading are then

$$(\sigma_x)_{y=+\frac{H}{2}} = (-0.50) \times \frac{2,500}{15} + (-2.74) \\ \times \frac{100,000}{15(30 \times 12)} = -83 - 51 = -134 \text{ lb per sq in.}$$

$$(\sigma_x)_{y=-\frac{H}{2}} = 1.24 \times \frac{2,500}{15} + 1.74 \\ \times \frac{100,000}{15(30 \times 12)} = 207 + 32 = +239 \text{ lb per sq in.}$$

Mention may be made here of another method for approximate determination of the stresses in a deep beam, suggested by A. P. Sinitsyn.<sup>10,11</sup> This method is based on the similarity between the differential equation for the Airy stress function and that for the deflection of a plate under transverse loading.

Use is made of this analogy, but, instead of solving for the analogous plate deflection from the differential equation, simplifying assumptions enable these deflections to be determined in an approximate manner that, essentially, neglects the torsional plate moments. The results undoubtedly are in error at certain places in which the assumptions made have a significant effect in changing the true

TABLE I.—COMPARISON OF RESULTS

Values of $H/L$	$(\sigma_x)_{y=+\frac{H}{2}}$		$(\sigma_x)_{y=-\frac{H}{2}}$	
	Sinitsyn's solution	Finite difference method	Sinitsyn's solution	Finite difference method
(a) UNIFORM LOAD $p$ AT BOTTOM EDGE				
$\frac{1}{2}$	-2.38	-1.82	+2.69	+2.71
1	-0.53	-0.50	+1.77	+1.24
(b) UNIFORM LOAD $p$ AT BOTTOM EDGE				
$\frac{1}{2}$	-2.33	-2.26	+2.62	+2.65
1	-0.32	-0.24	+1.52	+1.53

picture of the plate action. The accuracy of the results depends on the position of the points selected, the ratio of  $H/L$ , and the type of loading. This may be seen from the following comparison between the results of Sinitsyn's method and those given by the finite difference method. The discrepancy is substantial in certain cases. Stresses in Table 1 are in units of  $p/b$  and the midspan section of a beam is considered.

#### APPLICATION TO THE DESIGN OF REINFORCED CONCRETE STRUCTURES

For beams of homogeneous material, a determination of the bending and shear stresses by means of Figs. 2 to 6 gives sufficient information for design purposes. Reinforced concrete, on the other hand, is not a homogeneous material, so that one of the basic assumptions of the preceding data is not satisfied. The additional fact that the tension zone of reinforced concrete beams must be considered cracked, modifies the stress distribution even in shallow beams

<sup>10</sup> "Approximate Analysis of Beam-Walls" (in Russian), by A. P. Sinitsyn, *Proyekt y Standart*, No. 5, 1935, p. 21.

<sup>11</sup> *Ibid.*, N. 10, p. 24.

(Fig. 7). The analysis is simplified for such beams, however, by the reasonable validity of the assumption that plane cross sections remain plane despite cracking. For deep beams, on the other hand, this assumption cannot be made. The stress distribution in such reinforced concrete members must be expected to

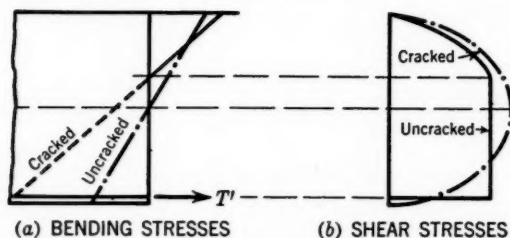


FIG. 7.—EFFECT OF CRACKING ON STRESS DISTRIBUTIONS IN A SHALLOW BEAM

differ from that given by Figs. 2 to 6 on two counts: (1) The nonhomogeneity of material; and (2) the cracking of the tension zone. A rigorous, theoretical analysis of the stresses in such beams is hardly feasible. For this reason, no unique and strictly justified design procedure can be proposed. Some dis-

cussion of this problem and suggestions regarding tension and shear reinforcement are given in the following sections.

**Tensile Reinforcement.**—The Portland Cement Association,<sup>5</sup> following European practice, suggests that the required steel area  $A_s$  be determined from

$$A_s = \frac{T}{f_s} \dots \dots \dots (4)$$

in which  $T$  is the total tension force computed from the homogeneous beam analysis (that is, Figs. 2 to 6) and  $f_s$  is the allowable unit stress of the steel. The entire tension reinforcement is located near the tension (bottom) edge. It must be realized that this procedure, although possibly safe in regard to strength, destroys completely the very assumptions on which the analysis is based and from which the value of  $T$  is determined. Indeed, by concentrating the tension reinforcement near the edge, the total tension force in the cracked section is forced to shift from the centroid of the tension areas of Figs. 2 to 6 to the centroid of the reinforcement (Fig. 8(a)). This is likely to increase the

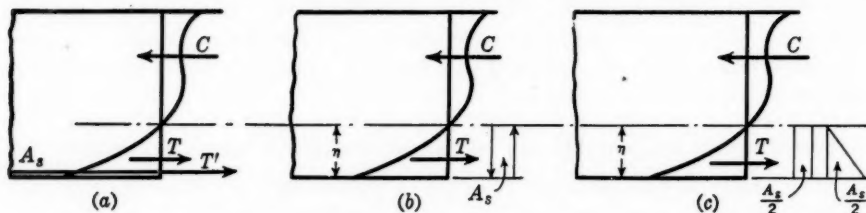


FIG. 8.—PROPOSED ARRANGEMENTS OF TENSILE REINFORCEMENT

effective internal lever arm although such a statement cannot be made with certainty for lack of rigorous, analytical support. If such increase does occur, it implies, of course, a decrease of the actual force  $T'$  in the reinforcement as compared to its computed value  $T$  and a corresponding decrease of the total compression force, since the couple formed by these forces equals the given



static moment. As a consequence, the resulting distribution of both normal and shear stresses is likely to be entirely different from that obtained from the analysis of the homogeneous beam. Even if a change of lever arm would not occur, the stress distribution would still be strongly affected by the forced shift of the tensile force to the edge since, in this case, the centroid of the compression area would have to shift by the same amount in the same direction to maintain equilibrium, with a consequent radical change in the distribution of the compression and shear stresses.

A further, practical argument against locating the reinforcement in this manner concerns the type and extent of cracking likely to occur. Cracks that begin at the tension edge would then be intercepted by reinforcement only near the edge, without any further reinforcement counteracting their growth in length and width up to the neutral axis. Since this axis, even in the homogeneous beam, is located at a distance from the tension edge varying from one half to one ninth of the depth—and likely to be larger yet in the nonhomogeneous, cracked beam—cracks of very considerable extent and width result in large, wall-like beams. To reduce these cracks additional distributed reinforcement at least of the amount usually prescribed against shrinkage would be called for over the entire tension zone.

Another scheme of reinforcing that could be rationally justified is one that would convert the beam as nearly as possible into a homogeneous beam. This can be achieved by distributing horizontal reinforcement uniformly throughout the tension zone to provide equal resistance and rigidity for every unit of depth throughout that zone, which is the essential prerequisite of "nearly homogeneous action." It is obvious that such distributed reinforcement would minimize cracking and would result in a location of the resultant of the steel stresses very close to, if not coincident with, that of the theoretical location of  $T$  (Fig. 8(b)).

To determine the amount of reinforcement required in this case, use can be made of the fact that from Figs. 2 to 6 the tensile stresses are seen to be nearly linearly distributed. Hence, with uniform steel distribution, the average steel stress for all practical purposes equals half that in the bar closest to the edge. The latter, as usual, must be limited to the allowable steel stress  $f_s$ . Hence

$$A_s = \frac{T}{f_{av}} = \frac{T}{0.5 f_s} = \frac{2 T}{f_s} \dots \dots \dots (5)$$

Comparison with Eq. 4 indicates that, for this manner of reinforcement, exactly twice as much steel would be required as in the first case, in which the entire steel was concentrated near the tension edge. It is likely, therefore, that such uniform reinforcement is overconservative and uneconomical, since the part of the steel near the neutral axis is stressed very little.

For an apparently sensible compromise between these conflicting requirements a different approach suggests itself.

The total required steel area may be determined from

$$A_s = \frac{1.5 T}{f_s} \dots \dots \dots (6)$$

Half of this area is uniformly distributed throughout the tension zone, while for the other half the spacing between rods is increased linearly with increasing distance from the tension edge, as schematically indicated in Fig. 8(c). Such distribution evidently provides sufficient steel at all levels to counteract excessive cracking and to secure at least some measure of homogeneous action. It can also be shown that the outermost rods of this arrangement will not be overstressed. Indeed, the arrangement of Fig. 8(b), which provides for the

TABLE 2—TOTAL TENSILE FORCE  $T$  AND DISTANCE  $\eta$  OF THE NEUTRAL AXIS FROM THE TENSION EDGE FOR SECTION  $x=0$

Types of loading	$\frac{H}{L}$	$T$	$\frac{\eta}{H}$
Uniform load $p$ at top edge (Fig. 1(a))	$\frac{1}{2}$	$0.510 p H$	0.400
	1	$0.131 p H$	0.255
	2	$0.010 p H$ $0.045 p H$	$\begin{Bmatrix} 0.890 \\ 0.583 \\ 0.111 \end{Bmatrix}$
Concentrated load $P$ at top edge (Fig. 1(b))	$\frac{1}{2}$	$1.125 P H/L$	0.547
	1	$\begin{Bmatrix} 0.022 P H/L \\ 0.196 P H/L \end{Bmatrix}$	$\begin{Bmatrix} 0.820 \\ 0.495 \\ 0.336 \end{Bmatrix}$
	2	$\begin{Bmatrix} 0.085 P H/L \\ 0.047 P H/L \end{Bmatrix}$	$\begin{Bmatrix} 0.900 \\ 0.505 \\ 0.117 \end{Bmatrix}$
Uniform load $p$ at bottom edge (Fig. 1(c))	$\frac{1}{2}$	$0.549 p H$	0.467
	1	$0.161 p H$	0.228
	2	$0.060 p H$	0.125
Concentrated load $P$ at bottom edge (Fig. 1(d))	$\frac{1}{2}$	$1.097 P H/L$	0.418
	1	$0.353 P H/L$	0.186
	2	$0.146 P H/L$	0.115
Two loads $P$ at third points of top edge (Fig. 1(e))	1	$0.359 P H/L$	0.278

$$a_s = \frac{2 T}{f_s \eta} \dots \dots \dots (7)$$

in which  $\eta$  is the distance to the neutral axis from the tension edge (see Fig. 8). For the arrangement of Fig. 8(c) and Eq. 6, the amount of reinforcement per vertical inch at the tension edge, where tensile stresses are highest, is easily computed from the separate contributions of the two total areas  $\frac{A_s}{2}$  each, and equals

$$a_s = \frac{0.75 T}{f_s \eta} + \frac{2 \times 0.75 T}{f_s \eta} = \frac{2.25 T}{f_s \eta} \dots (8)$$

Comparison of Eqs. 7 and 8 shows that this arrangement provides a steel density 12.5% larger than that required in scheme Fig. 8(b) to insure the edge steel against overstressing. (If this suggested scheme is followed, obviously no

actual distinction need be made between the rods of the two differently distributed areas  $\frac{A_s}{2}$ . Since, at the neutral axis,  $a_s = \frac{0.75 T}{f_s \eta}$ , comparison with Eq. 8

shows that such a steel distribution is achieved if the total area  $A_s$  is so arranged that the area per vertical inch or foot at the tension edge is three times that at the neutral axis, with reasonably linear transition between these two extremes.) It is not likely that the steel area required by this alternative is significantly larger than that for Eq. 4 and for Fig. 8(a), if, for the latter arrangement, the required additional shrinkage steel is added to the computed reinforcement  $A_s$ .

Since the use of any of these formulas for  $A_s$  requires that the magnitude of the total tensile force  $T$  and the location of the neutral axis be known, these values, measured from the stress curves, are given in Table 2.

It is seen from the stress distribution graphs that in some cases a second tensile region exists in addition to the bottom one. The stresses in these regions are often too small to cause cracks, making reinforcement unnecessary. If these stresses exceed the tensile strength of concrete, the steel area required can be easily computed from the corresponding value of  $T$ .

*Shear Reinforcement.*—The question of shear reinforcement is rather involved, even for shallow beams. The conventional procedure of designing for shear reinforcement is not theoretically exact, since it is merely an approximate means of accounting for the inclined principal tensile stresses in the concrete. In fact, beams do not fail by direct shear but by inclined tension induced by shear and normal stresses. The allowable shear stress  $v_c$  for concrete is determined empirically in such a manner that the value is low enough to insure against failure by inclined tension. Therefore, the allowable values of  $v_c$  are known to be considerably below the safe working stress of concrete in direct shear.<sup>12</sup>

For shallow, homogeneous beams in which the effect of  $\sigma_y$  is customarily neglected, the principal tensile stress  $\sigma_1$  is expressed by

$$\sigma_1 = \frac{\sigma_x}{2} + \frac{1}{2} \sqrt{\sigma_x^2 + 4\tau_{xy}^2} \dots \dots \dots (9)$$

and the angle of inclination  $\theta$  between  $\sigma_1$  and  $\sigma_x$  is given by

$$\tan 2\theta = \frac{2\tau_{xy}}{\sigma_x} \dots \dots \dots (10)$$

At the neutral axis  $\sigma_x = 0$  so that  $\sigma_1 = \tau_{xy}$  and  $\theta = 45^\circ$ . The same condition is assumed to be true in a reinforced concrete beam for any point below the neutral axis since, in view of cracking,  $\sigma_x$  is assumed to be zero throughout the tension part of the beam. Thus, the maximum inclined tension is, with sufficient accuracy, equal to the maximum shear stress, and web reinforcement is designed to resist the difference between this maximum shear stress and the amount of shear assigned to the concrete. It is well known that this method is a somewhat crude approximation and even slightly contradictory in itself, but it leads to a workable and safe design method for shallow beams.

However, even for shallow beams, Mr. Bay<sup>13,14,15</sup> has shown that the customary shear investigation leads to an overestimate of the magnitude of the maximum principal tensile stress and that the critical section for inclined tension is not at the support but at a distance of approximately  $0.65 H$  from the support. This is true since, near a support, the values of the vertical compressive stresses  $\sigma_y$  are large and cannot be neglected. Therefore, Eqs. 9 and

<sup>12</sup> "Design of Concrete Structures," by L. C. Urquhart and C. E. O'Rourke, McGraw-Hill Book Co., Inc., New York, N. Y., 4th Ed., 1940, p. 104.

<sup>13</sup> "Die schiefen Hauptzugspannungen beim Eisenbetonbalken," by H. Bay, *Ingenieur-Archiv*, Vol. 4, 1933, p. 244.

<sup>14</sup> "Scherbeanspruchung und Scherfestigkeit beim Beton," by H. Bay, *ibid.*, Vol. 14, 1943, pp. 267-276.

<sup>15</sup> "Der Einfluss der lotrechten Pressungen auf die Hauptzugspannungen beim Eisenbetonträger," by H. Bay, *Beton und Eisen*, Vol. 22, 1933, pp. 239-241.

10 are replaced by

$$\sigma_1 = \frac{\sigma_x + \sigma_y}{2} + \frac{1}{2} \sqrt{(\sigma_x - \sigma_y)^2 + 4\tau_{xy}^2} \dots \dots \dots (11)$$

$$\tan 2\theta = \frac{2\tau_{xy}}{\sigma_x - \sigma_y} \dots \dots \dots (12)$$

As seen from these equations, the presence of the compressive stress  $\sigma_y$  in the vicinity of the supports will reduce the magnitude of  $\sigma_1$  and also make  $\theta$  considerably less than  $45^\circ$ .

The principal tensile stresses are, therefore, nearly horizontal and may be carried by the longitudinal steel. Mr. Bay confirmed his analysis by tests that showed that the vicinity of the supports was free from cracks at loads far greater than those that had caused inclined cracks to occur in the same beams at points farther removed from the supports. The value of  $\sigma_y$  decreases with increasing distance from the support, so its effect on  $\sigma_1$  becomes negligible at a distance of about  $0.65 H$  from the support. Hence,  $\sigma_1$  reaches a maximum at this point. The tests by Mr. Bay also showed that the first inclined crack actually appeared at this place at which  $\sigma_1$  was a maximum.

The foregoing discussion is presented for the purpose of illustrating some of the difficulties involved in obtaining a logically consistent procedure for designing shear reinforcement. It is obvious that the conventional, semi-empirical method for shallow beams cannot be applied to deep beams in which the magnitude of  $\sigma_1$  is greatly affected by the presence of  $\sigma_y$  throughout the beam, in a manner similar to that just discussed for the vicinity of the supports of shallow beams. For these reasons the shear stress cannot be regarded as a reliable measure of inclined tension in deep beams, and accurate information could be obtained only by actually computing principal tension stresses. Since such a procedure would be extremely involved, the approximate method for shear in deep beams, proposed by the Portland Cement Association,<sup>5</sup> appears sensible. It is suggested that, for  $H/L > 0.4$ , the allowable shear stress for deep beams be

increased to  $\frac{v_c \left(1 + 5 \frac{H}{L}\right)}{3}$  in which  $v_c$  is the value allowed for shallow beams.

This accounts for the fact that, because of the presence of vertical compression  $\sigma_y$ , the inclined tension  $\sigma_1$  is smaller than the maximum shear  $\tau_{\max}$ . Instead of computing this lower inclined tension and comparing it with  $v_c$ , the Portland Cement Association method proposes raising the allowable shear stress and comparing it with  $\tau_{\max}$ . The effect is identical, of course, at least qualitatively.

It is further recommended in the previously mentioned publication that the maximum shear be computed from  $\tau_{\max} = \frac{8V}{7bH}$ , that is, from the usual, approximate formula for shallow beams. This procedure may be justifiable for lack of better information if the beam is reinforced according to Fig. 8(a) (with reinforcement at the tension edge) since, in that case, the actual stress distribution is highly uncertain. However, if reinforcement is arranged to approach more closely the condition of homogeneity (such as in Fig. 8(b) or 8(c)), the shear distribution can be assumed to be fairly close to that in a homogeneous beam, and  $\tau_{\max}$  can then be taken from the appropriate curves of Figs. 2 to 6.

In many cases  $\tau_{\max}$  will be found to be smaller than the shear value assigned to concrete in deep beams, so that shear reinforcement is not required. If this is not the case such reinforcement must be provided primarily by inclined bars or by a network of horizontal and vertical bars. In contrast to shallow beams, cracks in deep beams are nearly vertical (since principal tensions are nearly horizontal) so that vertical shear reinforcement, such as stirrups, is largely ineffective.

This brief discussion of the questions involved in correctly reinforcing deep concrete beams is not meant to be either exhaustive or authoritative. It is merely intended to point up some of the differences in approach required in designing deep as compared to shallow concrete beams and to suggest methods that seem reasonable to the authors. Experimental investigations in this field would do more to settle the question of the most desirable arrangement of reinforcement than any amount of theorizing.

#### CONCLUSIONS

For shallow beams, the normal pressures on the longitudinal edges have little effect on the stress distributions at sections a small distance from supports or concentrated loads. The stresses at these sections depend only on the bending moment and shear force, and the values computed by the ordinary beam formulas are accurate for practical purposes. The stresses in deep beams, however, depend not only on the bending moment and shear force at a section but also on the variation of normal pressures along the loaded edges. The greater the height-to-span ratio, the more significant does the latter factor become. Thus, for deep beams in which the load is applied to the bottom edge, the stress distributions are quite different from those caused by the same loading applied at the top edge. In general, the ordinary beam formulas may be considered as valid for computing stresses in beams whose height-to-span ratio is less than one half and in which the sections under study are not too near to loads and reactions.

Analytical results for the stresses in deep beams of homogeneous material are presented and can be applied directly to the design of structures made of such material. The design of deep reinforced concrete beams, on the other hand, involves the difficulty of dealing with a slightly cracked, nonhomogeneous material. A rigorous theoretical analysis of this situation is hardly possible. The stresses obtained from analyses based on the assumption of homogeneity may serve as a reasonable guide for estimating the actual stresses in the cracked state of the nonhomogeneous beam, provided the beam is so reinforced as to approach homogeneous performance. Suggestions regarding tensile and shear reinforcement have been presented to achieve such performance and to limit the extent of cracking.

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